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A GROUND-MOTION DISPLACEMENT GAUGE, (U)

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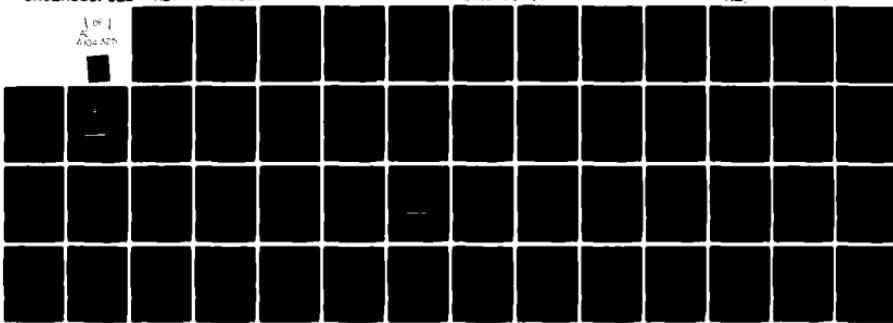
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## A GROUND-MOTION DISPLACEMENT GAUGE

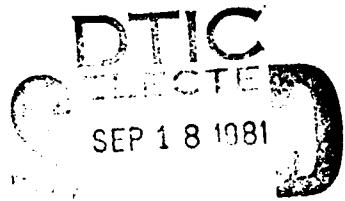
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1 August 1980

Topical Report for Period 5 January 1980—1 August 1980

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) This report describes a ground-motion displacement gauge and its applications. The gauge concept is based upon the free-fall time of a mass released inside a gauge cannister prior to the onset of ground motion. The time-of-fall of the mass to the cannister bottom during ground displacement, together with the free-fall time for the gauge at rest, uniquely determines the vertical displacement at one time. Simulation of ground displacement by gauge drop tests demonstrated the operating principle.		

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## SUMMARY

This report summarizes the concept, applications, and results of developmental tests for a special ground-motion displacement gauge. Other possible uses, which have not been tested, are also discussed.

The gauge concept is based upon the free-fall time of a mass released inside a gauge cannister prior to the onset of ground motion. The time-of-fall of the mass to the cannister bottom during ground displacement, together with the free-fall time for the gauge at rest, uniquely determines the vertical displacement at one time. Simulation of ground displacement by gauge drop tests demonstrated the operating principle.

One application of the gauge, for which it was originally proposed, would be to measure permanent displacements produced by underground explosions. The displacement potential determined from the measurement would provide a source strength of seismic signals from the explosion. Another use is to verify the measurements of ground displacement obtained by integration of particle velocity or acceleration gauge data. The size of the gauge is one limiting factor on the range of displacement which can be measured. Another is the ground stress which the gauge must survive. The practical range of measurements would be from .1- to 100-cm displacement.

## PREFACE

This report describes a ground-motion displacement gauge previously proposed by RDA. Publication of the report was prompted by suggestions from Mr. R. J. Port of RDA for new applications in high-explosive simulation programs. Dr. E. A. Martinelli offered valuable assistance on other applications discussed in the report.

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## I. INTRODUCTION

This report describes the concept and development of a ground-motion displacement gauge suggested by the author at RDA in 1972 (Ref. 1). During the period 1972 through 1973, the MIGHTY MITE series of nuclear and high-explosive underground tests were conducted at the Nevada Test Site. The objective of these tests was to determine the effectiveness of various heat sink concepts for quenching explosions inside cavities. The Advanced Research Projects Agency (ARPA) sponsored the test program which was executed by the Defense Nuclear Agency. ARPA was primarily concerned that, in the event of complete test ban negotiations, clandestine tests in underground cavities could deny seismic detection. Reference 2 describes these tests.

From the theory of cavity decoupling, it has been shown that the seismic energy transmitted to teleseismic distances ( $\sim 10^3$  km) is related to the permanent elastic particle displacement at or near the cavity wall (Ref. 3).

In Appendix A a derivation is given for the relation between permanent displacement and the energy released in a cavity. This application of a permanent ground-motion displacement measurement motivated the concept of the gauge discussed in this report.

## II. DESCRIPTION OF THE VERTICAL DISPLACEMENT GAUGE

### 1. OPERATING PRINCIPLE

The principle of gauge operation is illustrated using Figures 1 and 2. In Figure 1, the initial position of the gauge case is shown in the solid figure. Any point A of the gauge case moves to a new position B along the displacement path. At time prior to the beginning of motion, a mass is released at the top of the gauge. On impact of the mass shorting pin upon the foil switch located at the bottom of the gauge case, the time of impact is recorded with the gauge now at position B. During the time-of-fall, the mass drops through the distance  $\ell$  and the gauge is displaced vertically by the amount  $d$  as shown. For the case of no vertical motion, the mass would drop a fixed distance  $S$  as shown in Figure 1. From the recorded release and impact times,  $t_r$  and  $t_i$  respectively, the displacement is determined by

$$d = S - \frac{1}{2}g(t_i - t_r)^2$$

or alternatively,

$$d = \frac{1}{2}g \left\{ t_o^2 - (t_i - t_r)^2 \right\} \quad (1)$$

where

$$S = \frac{1}{2}gt_o^2 \quad (2)$$

and  $t_o$  is the drop time for zero displacement. If  $t_i - t_r > t_o$ , the displacement is downward. If  $t_i - t_r < t_o$ , the motion is upward.

It is emphasized that so far this gauge design measures only vertical displacement at just one point.

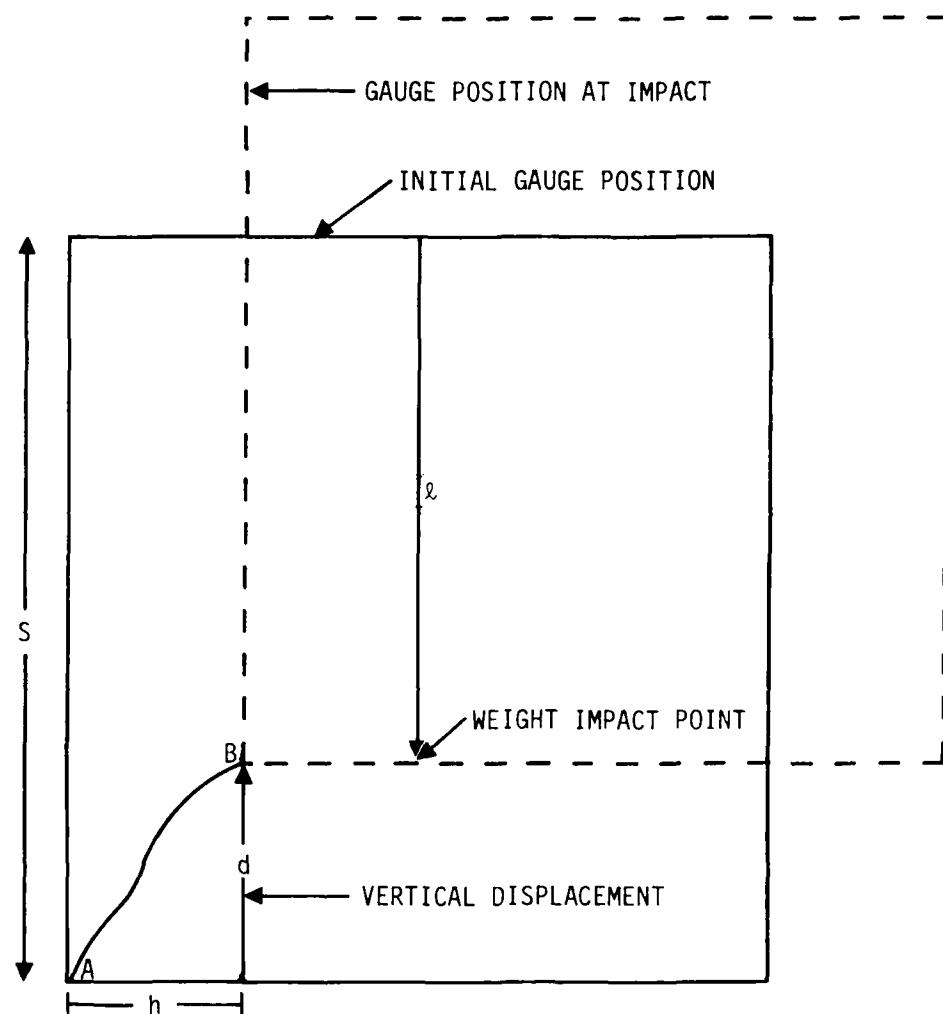


Figure 1. Principle of ground-motion gauge operation.

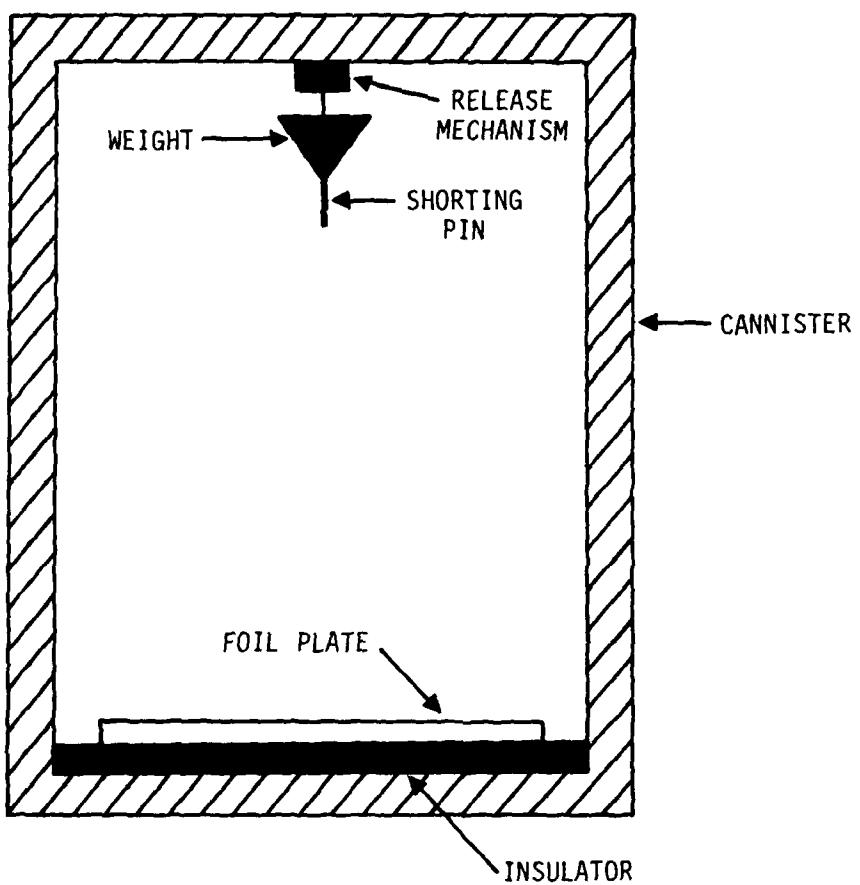


Figure 2. Schematic showing basic components of ground-motion displacement gauge.

## 2. OPERATIONAL CONSTRAINTS

Several features of this technique may limit the application of the gauge to a specific range of displacement measurements. Gauge construction must be preceded by careful estimates of the maximum motions and durations expected.

### a. Gauge dimensions

(1) Upward displacement--The maximum upward displacement which can be measured must be less than  $S$ , otherwise the weight impact would occur prior to the peak value of displacement. Thus the vertical dimensions must be greater than  $1/2gt_0^2$ . Figure 3 shows a plot of Eq. 2 to estimate the gauge size  $S$ . It is noted that the gauge size rapidly increases with the time-of-fall for upward motions.

In Figure 4, two upward-displacement waveforms are shown corresponding to a cube-root yield scaling factor of 2. The dashed curves are the free-fall paths for the different gauge sizes shown and for zero time of release. A mass drop height of almost 8 cm is required to measure displacement beyond the peak for  $W = 1$ , while about 4 cm is required for the smaller waveform. For smaller drop heights  $S$ , impact occurs prior to the peak displacement. It can also be noted that the gauge size  $S$  does not scale with  $W^{1/3}$  since the time-of-fall of the weight does not scale.

Another parameter which could be useful is controlled advance or delay time-of-weight release. Figure 5 illustrates these effects of advance or delay in the release time relative to zero time of release for upward displacement. The figure shows the different impact times for releases at -15 and +15 ms for 2.54-cm and 5-cm gauge sizes.

(2) Downward displacement--Downward displacements present different problems since the gauge motion is following the free-fall of the weight. The downward displacement

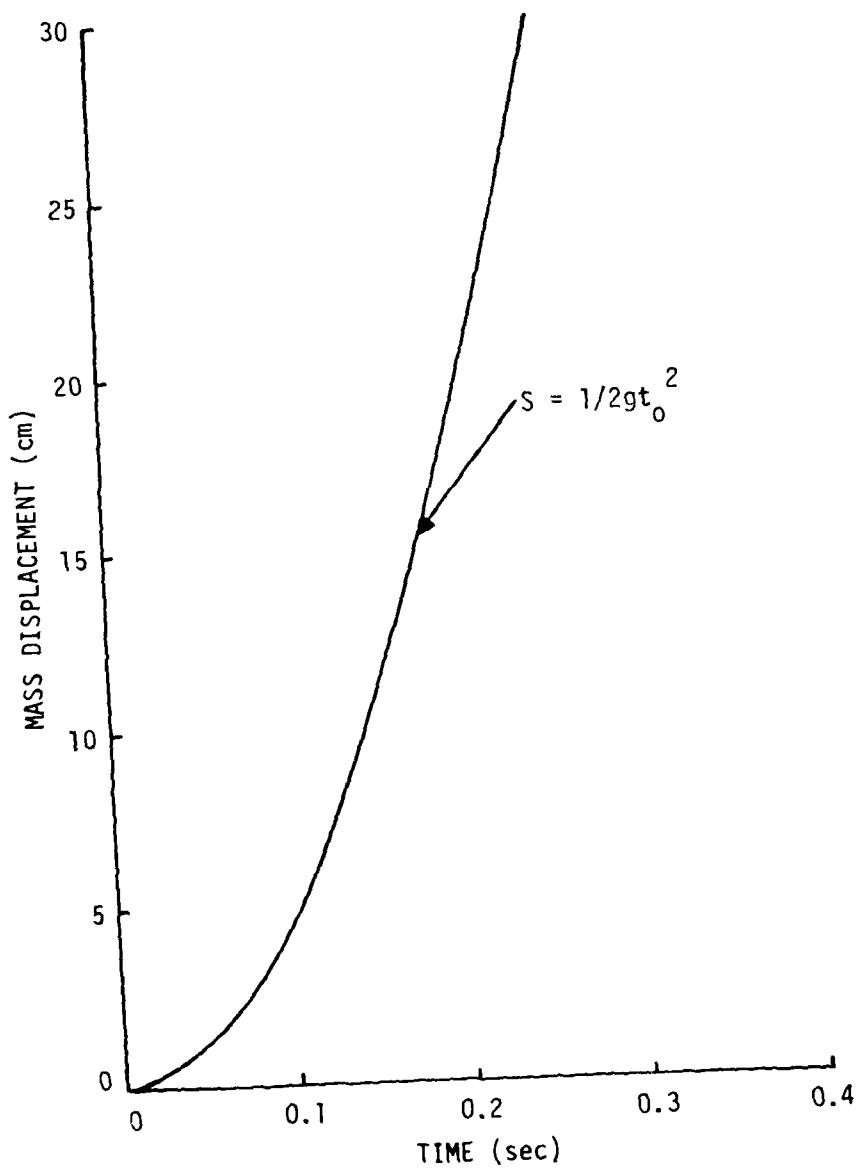


Figure 3. Estimate of gauge size vs zero-displacement drop time.

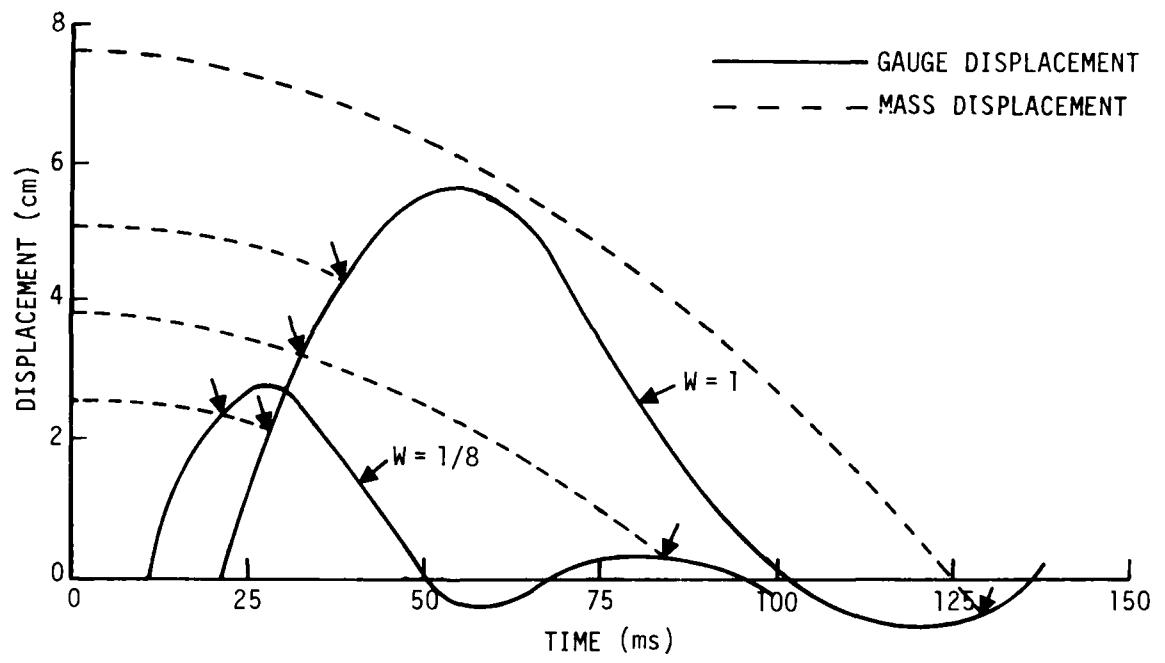


Figure 4. Illustration of displacement at impact time for varying gauge sizes and waveforms.

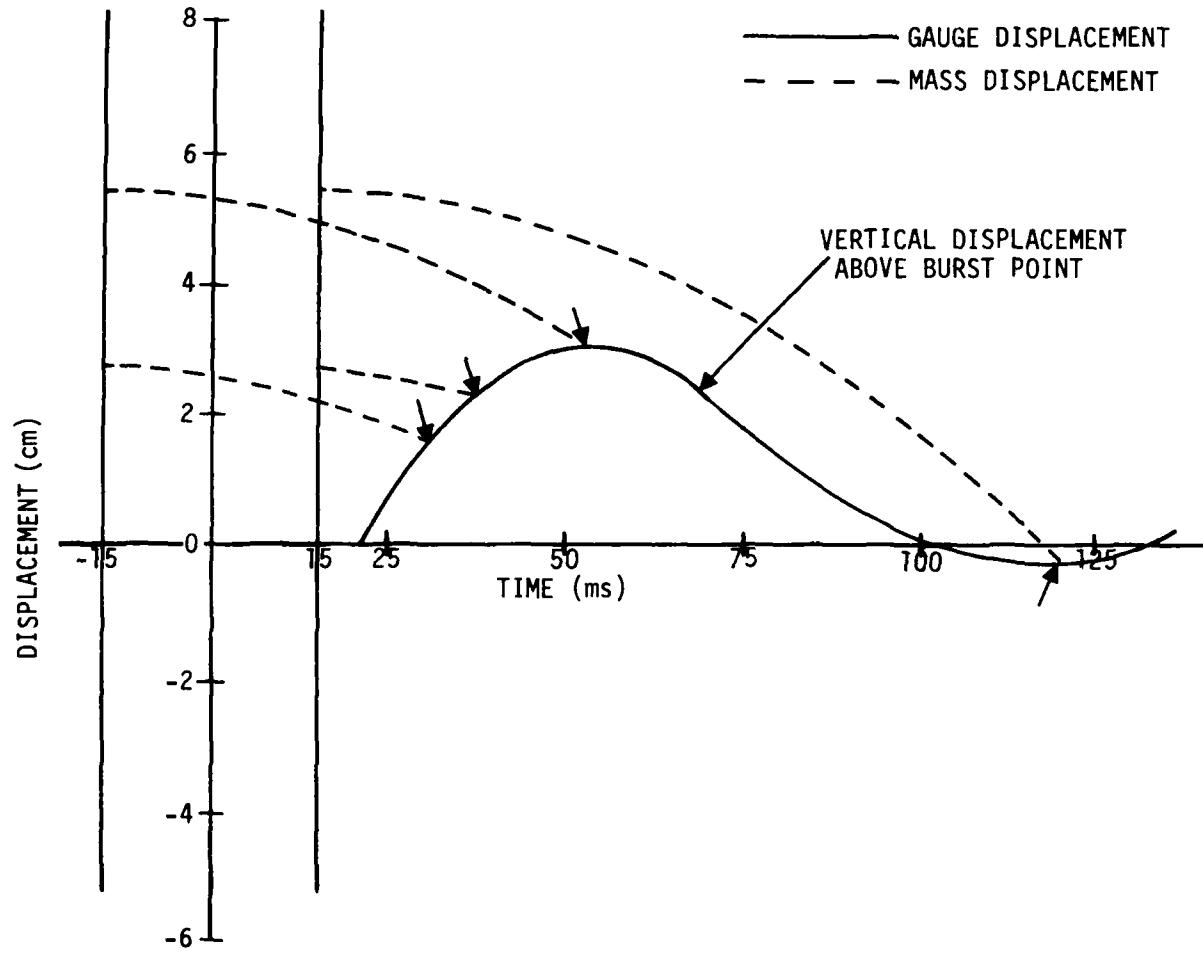


Figure 5. Effects of gauge size and weight-release time on recorded upward displacement.

$d(t)$  at any time  $t$  must not exceed the free-fall position of the mass; namely,

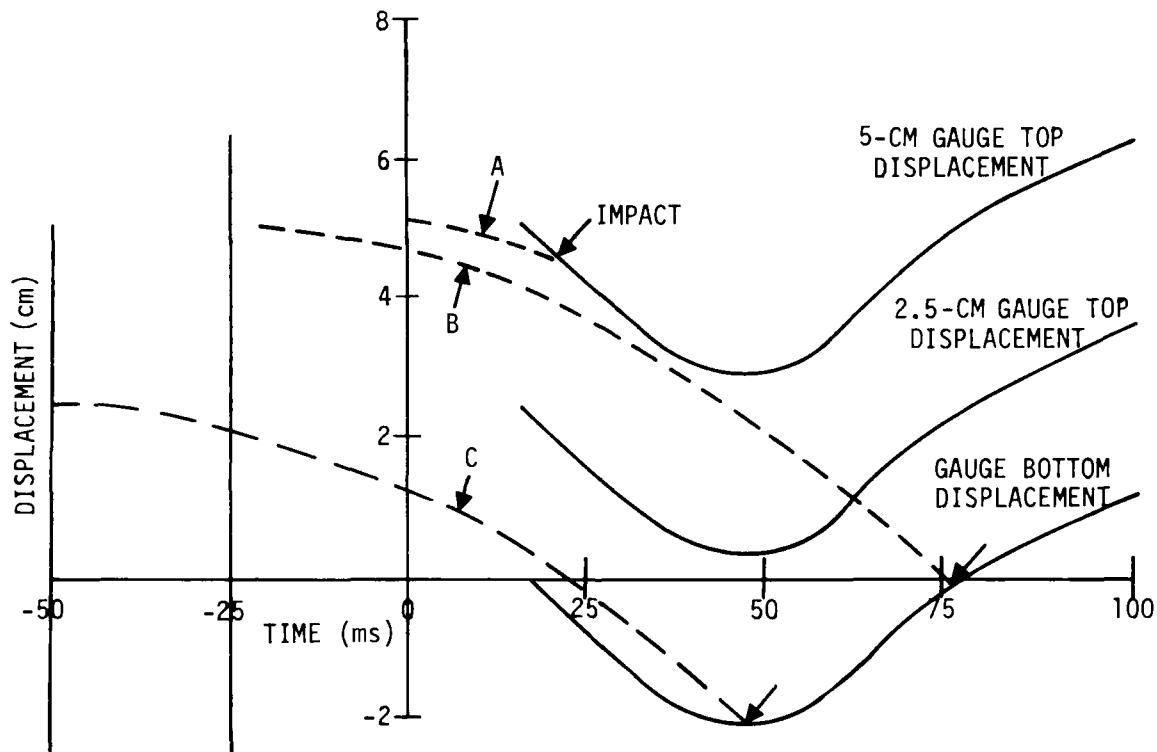
$$d(t) \leq 1/2g(t - t_r)^2$$

Otherwise, the top of the gauge case would overtake the falling weight.

Several different modes of gauge operation will affect the measured displacement for downward motion. Figure 6 illustrates these cases. Traces of the particle motion of the top and bottom points of the gauge are shown.

- Curve A illustrates a case where the gauge overtakes the mass released at zero time, which would result in failure of the measurement.
- Curve B shows the effect of pre-release of the mass at -25 ms and for the same initial mass height of 2.54 cm. In this case, the weight trajectory does not intersect that of the gauge top.
- Curve C shows the effect of reducing the gauge size from 5 cm to 2.5 cm and pre-release of the weight at -50 ms.

(3) Problem with horizontal motion--Referring to Figure 1, the horizontal gauge displacement  $h$  is shown at position B. From the figure it can be seen that the maximum horizontal displacement at impact must be less than the initial distance to the gauge wall. Otherwise, the mass would strike the wall and interfere with the time-of-fall measurement. The horizontal dimensions then must be taken into account in the specific application of the gauge. Figure 7 illustrates the complicated particle motion below the ground surface from a surface burst. The gauge would require more than 5 cm



**Figure 6.** Effects of mass-release time and gauge size upon measured displacement of a downward vertical motion.

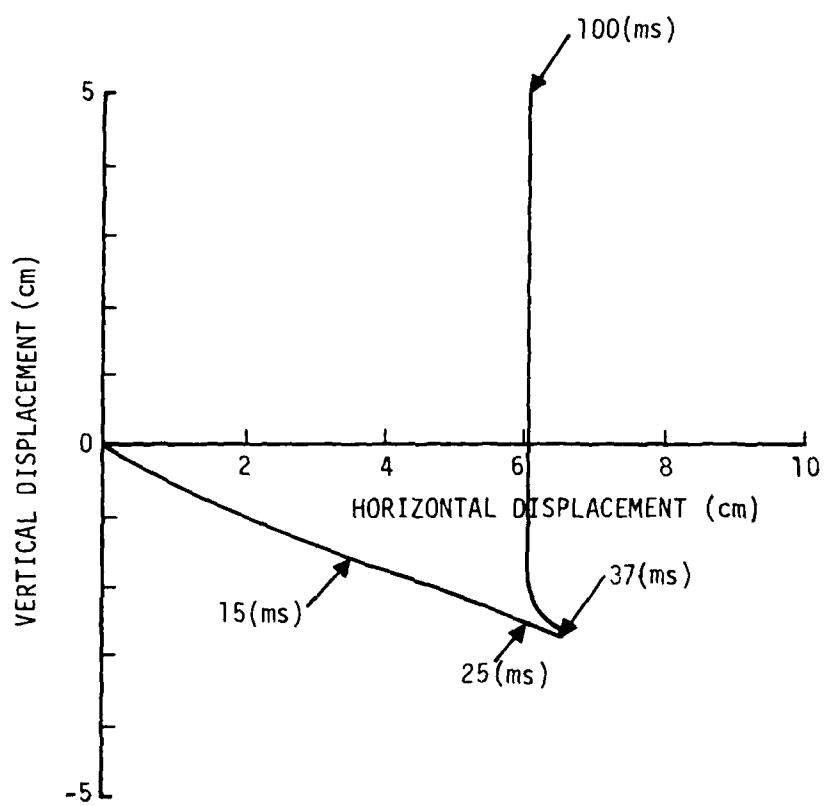


Figure 7. Particle displacement at 12-m range and 9-m depth below ground surface from a 500-ton surface burst.

horizontal clearance between the initial weight position and the gauge wall to record vertical motion to times greater than 37 ms.

(4) Gauge case strength--Mechanical hardening requirements for the displacement gauge may be greater than those for other types of ground-motion gauges. The large void required inside the gauge casing to accommodate the horizontal and vertical motions of the falling weight will increase the structural strength necessary to prevent buckling or crushing of the case prior to the measurement. The stress level at which a desired displacement measurement is to be made will depend upon the source yield and the peak displacement. All these factors must be taken into account in the design of the gauge housing. A great deal of experimental data on shock loading of cylindrical and other structures is available to estimate the structural hardening requirements or limitations for the design of the displacement gauge.

### 3. REVIEW OF ADVANTAGES AND DISADVANTAGES OF THE GAUGE DESIGN

These are some of the advantages associated with the gauge design described:

- No mechanical or electronic integration is required.
- There is an accurate inertial frame of reference (free fall).
- The data recorded includes only time-interval measurements, while amplitude measurements require calibration.
- Time-base drift is not possible.
- Time measurements of switch opening and closure provide high signal-to-noise ratio.

This gauge design does have certain disadvantages, such as:

- Only the vertical component of displacement is possible in the present gauge design.
- The dependence of gauge size upon magnitude of displacement requires careful estimates of expected ground motions.
- The time-squared dependence of free-fall distance imposes practical limitations on gauge size.

#### 4. SOURCES OF ERROR IN MEASUREMENTS

Two possible effects on the accuracy of measurement are considered.

- Air drag on drop mass. Relative motion of the drop mass and the air enclosed in the gauge canister will induce air drag, which will oppose the free-fall motion of the weight. In Appendix B, this problem is analyzed with estimates of errors expected.
- Rotational motion effects on measurements. If the gauge motion undergoes a rotation, the effect on the measurement is to introduce an error in the measurement. In Appendix C, this effect is discussed.

### III. DEVELOPMENT AND TESTS OF PROTOTYPE DISPLACEMENT GAUGE

#### 1. GAUGE CONSTRUCTION

A prototype displacement gauge was built and tested by Systems, Science and Software at the Green Farm Test Facility in March 1974. Specific design features of the gauge were contributed by the author and Dr. Howard Kratz (formerly of Systems, Science and Software). Electronic circuitry design and laboratory testing were performed by Mr. William Muhl of IMED Corporation (Ref. 4). Figure 8 is an engineering sketch of the full-scale model with specific features of the gauge indicated.

a. Mass-release mechanism--The drop mass is supported by the copper wire under tension from the leaf spring. An electrical circuit (normally closed) is completed through the breakwire to the drop mass and mass support. On receiving a release signal, the solenoid is energized, breaking the wire which opens the electrical circuit and at the same time releases the drop mass. Opening of the circuit supplies an electrical timing signal of the mass release.

b. Foil switch--The foil switch, located at the bottom of the gauge case, consists of a copper foil insulated from the metal case by a mylar foil. An electrical circuit (initially open at the switch) is completed by the shorting pin attached to the drop mass, generating an electrical timing signal at impact.

c. Electrical circuitry--A block diagram of the electrical circuitry is shown in Figure 9. The mass-release delay circuit generates a signal prior to shock arrival at the gauge. Timing of the mass release could occur at or prior to detonation time, or be delayed depending upon the distance between the gauge and explosive charge.

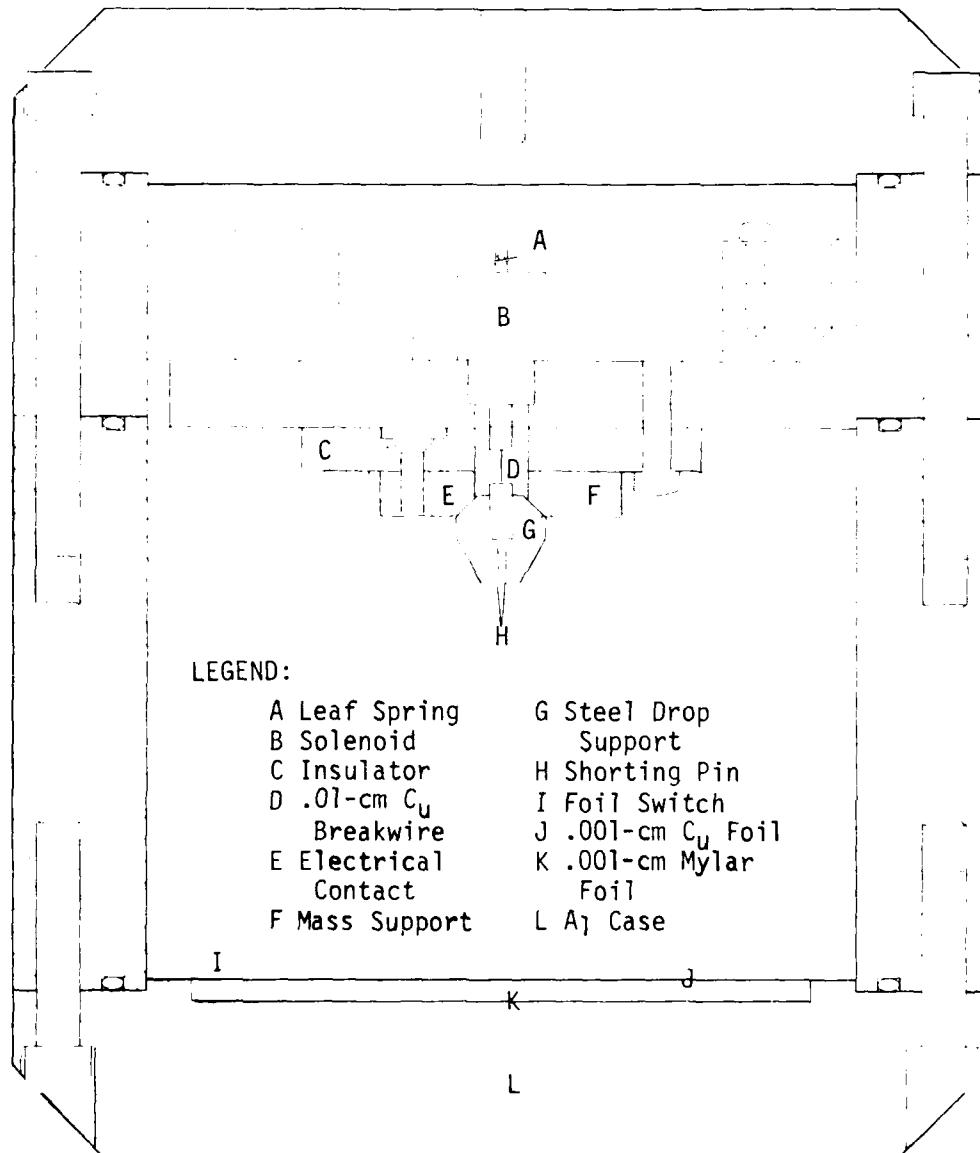


Figure 8. Full-scale drawing of prototype ground-motion displacement gauge.

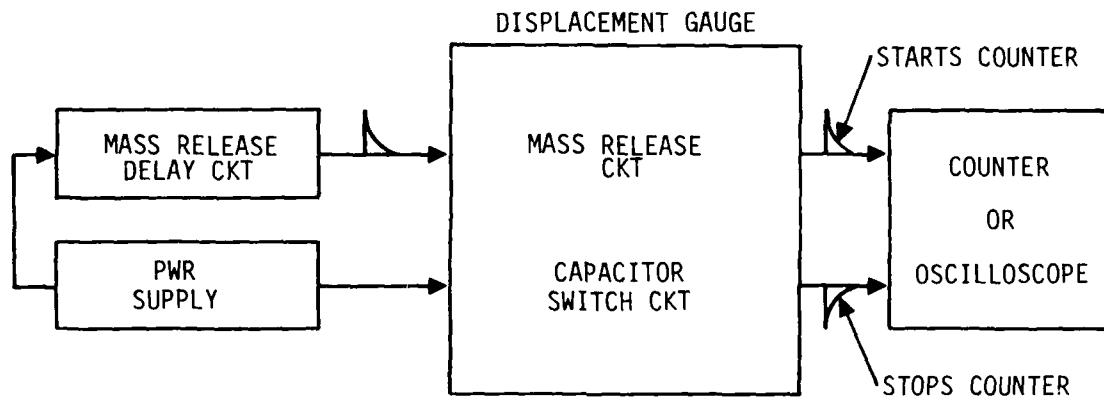


Figure 9. Block diagram of displacement gauge system.

The signal generated by the breakwire initiates the horizontal sweep of an oscilloscope, or can be used to start a counter. Arrival of the impact signal is recorded either on the oscilloscope or stops the counter if used. The time difference provides the time-of-fall from which the vertical displacement is calculated. Figure 10 is a schematic diagram showing the signal-generating circuits.

## 2. LABORATORY TESTS AND DISCUSSION OF RESULTS

During the MIGHTY MITE series of underground tests on seismic decoupling (Ref. 1), a test was planned that included the displacement gauge concept in the test program. That test was later cancelled; however, funding for construction and laboratory testing of the gauge was approved. Although no ground-motion displacement tests were attempted, gauge-displacement experiments were conducted to verify the gauge concept. Two sets of experiments are described and the results are presented in the following paragraphs.

a. Time-of-fall measurements--The first series of experiments simply consisted of measuring the free-fall time of the drop weight through an accurately measured drop distance with no motion of the gauge case. All the electronic circuits were energized as previously discussed. A trigger signal applied to the weight-release mechanism started an electronic counter with a time resolution of a few  $\mu$ sec. A signal at impact stopped the counter from which the time-of-fall was recorded. Table 1 summarizes these results obtained from the given drop distance. Also listed are the time differences from the calculated value, and the percentage differences. The consistency in the measurements would seem to indicate a systematic error either in the timing measurements or in the recorded drop distance. The difference between  $t_0$  and the mean is 2.0 ms.

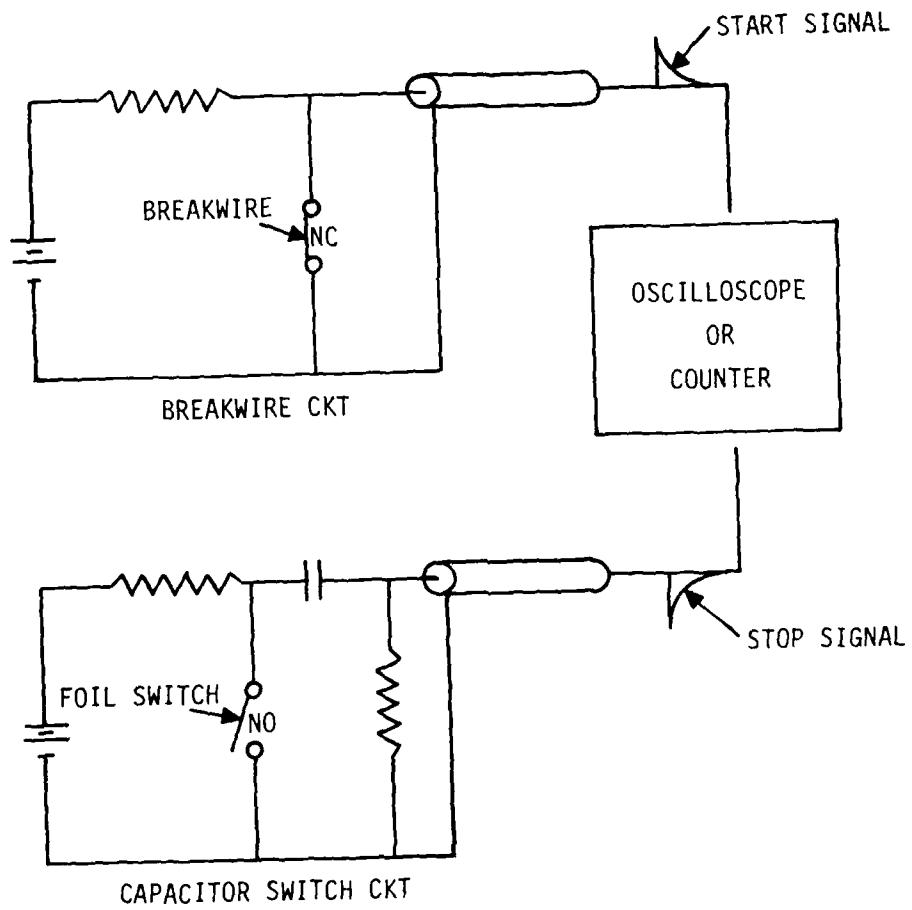


Figure 10. Schematic diagram of displacement gauge circuitry.

TABLE 1. FREE-FALL DROP-TEST MEASUREMENTS

Drop Distance  $S = 3.498 \text{ cm}$   
 Acceleration  $g = 980 \text{ cm/sec}^2$   
 Calculated Time-of-Fall  $t_0 = 84.486 \text{ ms}$

<u>TEST NO.</u>	<u>TIME-OF-FALL (ms)</u>	<u><math>(t_0 - t) \text{ms}</math></u>	<u>% DIFF</u>
1	82.52	1.967	2.32
2	82.64	1.845	2.18
3	82.21	2.276	2.69
4	82.70	1.786	2.11
5	82.69	1.796	2.13
6	82.60	1.886	2.23
7	82.22	2.266	2.68
8	82.09	2.396	2.84

$$\bar{t} = 82.46(\text{ms})$$

$$\sigma = 0.230(\text{ms})$$

b. Gauge drop tests--To simulate ground-motion displacement tests, the gauge was suspended a measured distance  $d_m$  above a bench top as shown in Figure 11. The gauge canister was released by firing of the HE detonator just after the drop mass was released. Measurements of time-of-fall  $t_m$  of the mass were recorded from which the apparent displacement was calculated by

$$d_a = s - \frac{1}{2}gt_m^2$$

where  $s$  is the initial pin distance from the bottom of the gauge.

Table 2 summarizes the test results. The apparent displacements are significantly different from the displacements (gauge drop distances) measured prior to each test by more than 10 percent. Again, as noted in Table 1, the time-of-fall in each case is less than was predicted from the known displacement. Applying the correction of 2 ms from Table 1, the corrected percentages of difference in the displacements are reduced as shown in the last column of Table 2. These results do not determine the source of the apparent systematic error, and time did not permit further investigation of the problem. There is little reason to doubt that the accuracy of the measurements can be significantly improved by troubleshooting the weight release and impact timing.

TABLE 2. DROP-TEST MEASUREMENTS

$$s = 5.08 \text{ cm}; \quad g = 980 \text{ cm/sec}^2$$

TEST NO.	$t_m$ (ms)	$d_m$ (cm)	$d_a$ (cm)	% DIFF	% DIFF (corr)
1	123.575	2.774	2.403	13.4	4.58
2	123.294	2.703	2.368	12.5	3.36
3	122.060	2.530	2.220	12.2	2.70
4	125.090	2.891	2.5873	10.5	1.95

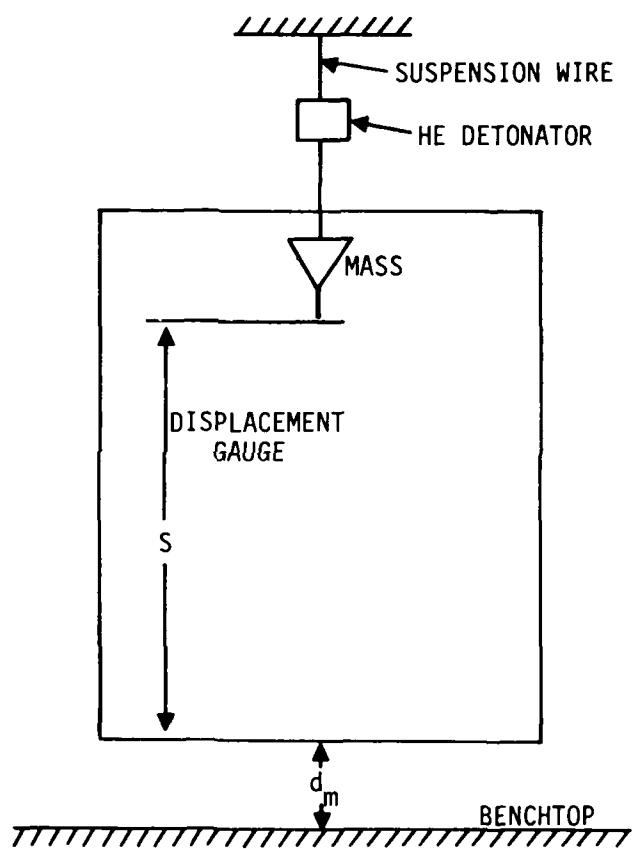


Figure 11. Experimental setup for simulation tests of displacement gauge.

#### IV. MODIFICATION OF GAUGE DESIGN FOR OTHER USES

##### 1. PARTICLE VELOCITY

Measurement of particle velocity, in addition to particle displacement, would provide another data check on existing particle velocity gauges. The same principle of time-interval measurements is all that would be required using the concept that will be described in this section.

Figure 12 illustrates the concept. Two separate displacement gauges are contained within the same cannister as shown. Two drop masses and two impact switches would provide two near-displacement measurements separated by a small time interval. The spacing of the two measurements could be accomplished in either of two ways. Shorting pins of unequal length released simultaneously would provide two displacements. From Figure 12 the masses released at  $t_r$  would impact at times  $t_1$  and  $t_2$ . The corresponding displacements are given by

$$d_1 = s_1 - 1/2g(t_1 - t_r)^2$$

$$d_2 = s_2 - 1/2g(t_2 - t_r)^2,$$

from which the average velocity over the displacement interval is obtained:

$$\bar{u} = \frac{d_2 - d_1}{t_2 - t_1} .$$

Figure 13 illustrates the principle of velocity measurement. Two different displacement curves are shown to illustrate positive and negative velocity measurements.

The second method works basically in the same manner where the two masses are released at different times,

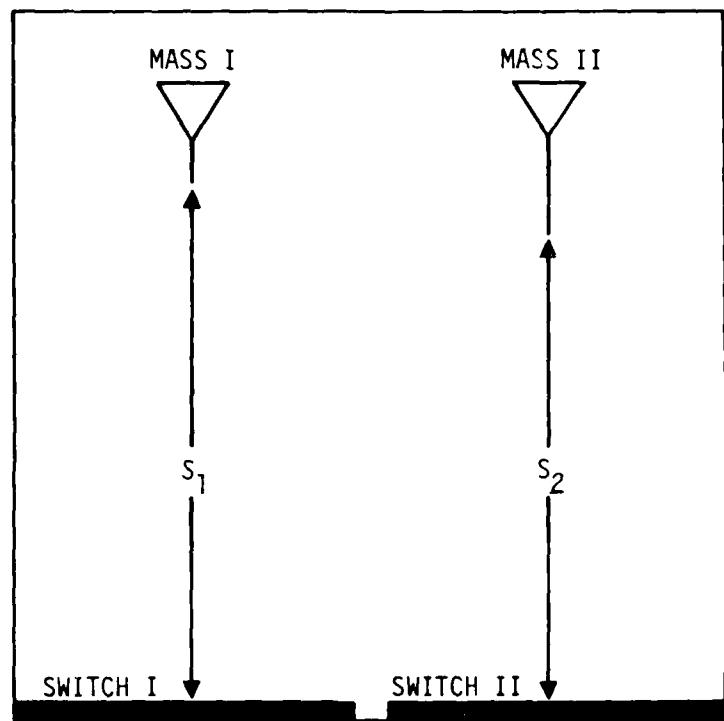


Figure 12. Modified displacement velocity gauge.

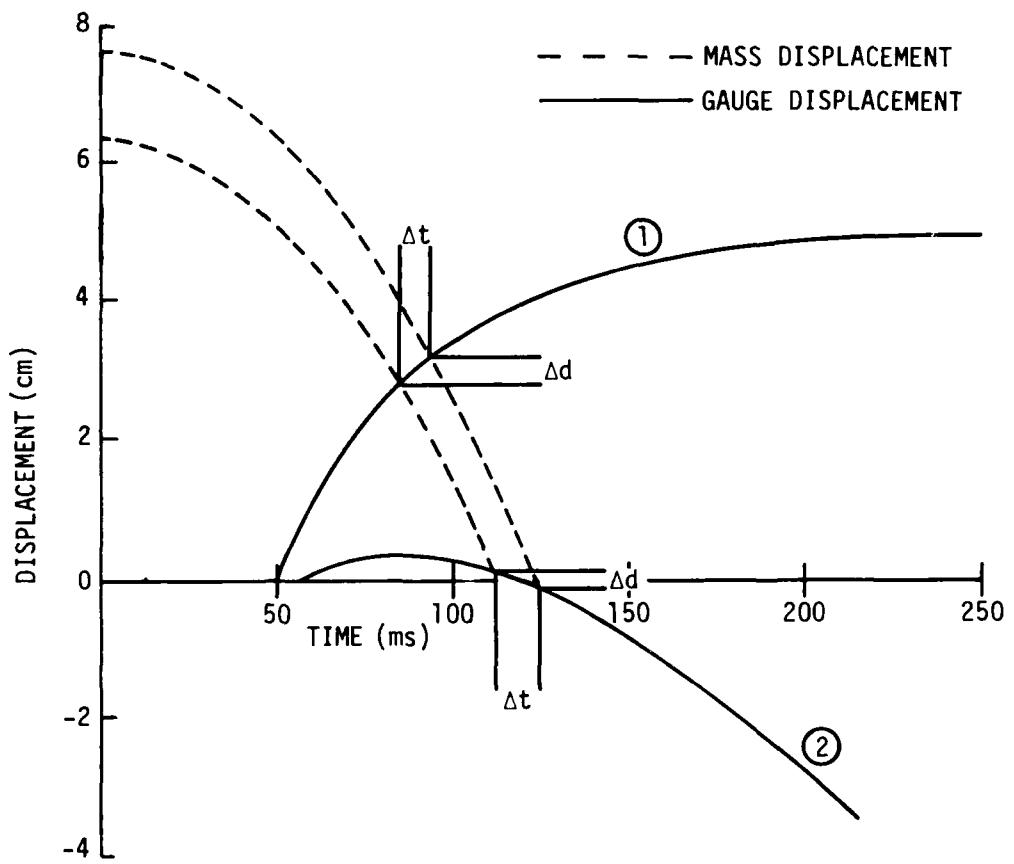


Figure 13. Particle velocity obtained from displacement measurements.

$t_1$  and  $t_2$ , with impact times  $t_3$  and  $t_4$ , respectively. The displacements are obtained from

$$d_1 = S - 1/2g(t_3 - t_1)^2$$

$$d_2 = S - 1/2g(t_4 - t_2)^2$$

and the velocity calculated from

$$\bar{u} = \frac{d_2 - d_1}{t_4 - t_3} .$$

## 2. HORIZONTAL AND VERTICAL DISPLACEMENT

Another possible application of the gauge concept is the measurement of the horizontal and vertical components of permanent displacement. Only the vertical component can be measured using the basic design as illustrated in Figure 1. The horizontal motion could be inferred from the vertical measurement in the case of known spherical symmetry along a given radial direction.

A gauge of special design, used in conjunction with the basic model, could in principle be used to determine both components of permanent displacement. Figure 14 is an illustration of the dual gauge principle where Gauge I has been modified by constructing the foil switch in the form of a cone as shown. Gauge II is the basic design discussed throughout the report. Operation of the system is as follows: at 'zero time' both weights are released at a distance  $S$  from the common base BB. The gauge undergoes a final upward displacement  $D$ , with horizontal and vertical components  $H$  and  $V$  as shown. Impacts of the masses at  $P_1$  and  $P_2$  occur at times  $t_1$  and  $t_2$  respectively, with the base now located at  $B'B'$ . Motions of the drop masses at impact are given by

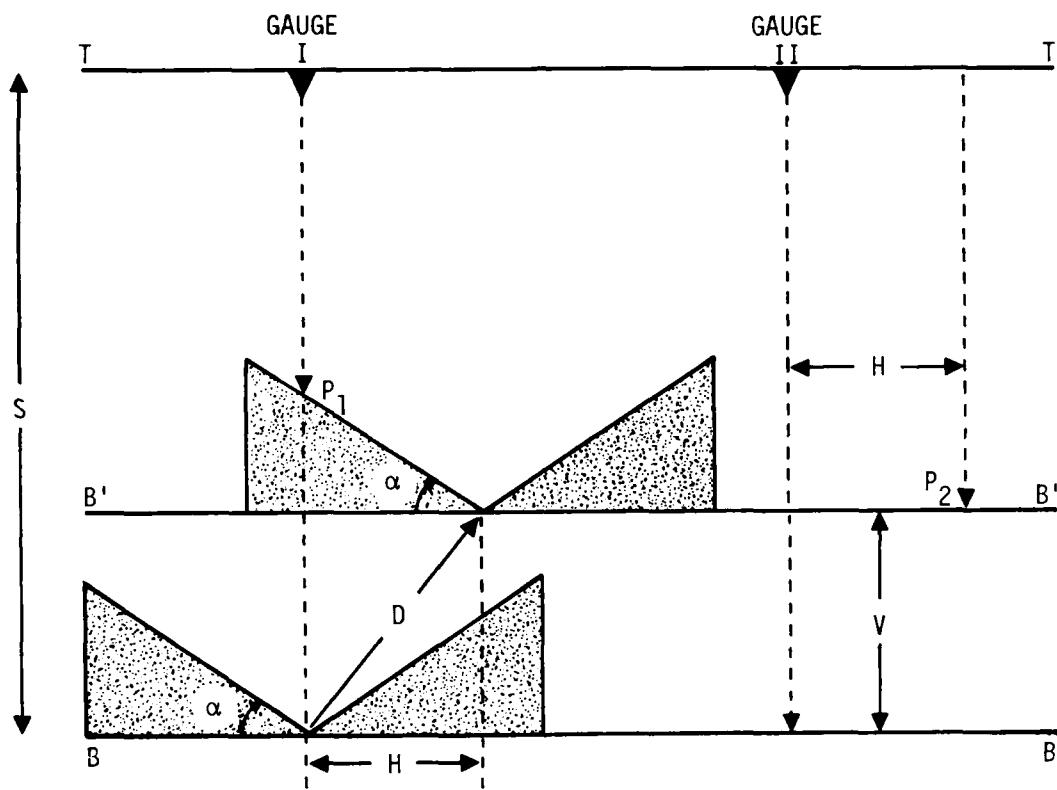


Figure 14. Schematic illustration of technique to measure horizontal and vertical displacement.

$$\frac{1}{2}gt_1^2 = s - v - H\tan\alpha$$

$$\frac{1}{2}gt_2^2 = s - v,$$

hence,

$$\frac{1}{2}g(t_2^2 - t_1^2) = H\tan\alpha,$$

where  $\alpha$  is the angle subtended from the horizontal by the conical foil switch. The displacement can then be calculated to be

$$H = \frac{1}{2}g(t_2^2 - t_1^2)/\tan\alpha$$

$$v = s - \frac{1}{2}gt_2^2$$

It should be noted that this technique can only be used for a final displacement since, in the case of a transient measurement, two different displacements would occur at times  $t_1$  and  $t_2$ , which would not uniquely determine these displacements.

## V. VERIFICATION TESTS OF ACCELEROMETER AND VELOCITY GAUGES

More recently, another application has been suggested. Particularly with the increase in the application of high explosives for the simulation of nuclear explosions, the gauge can be an important tool in calibrating ground-motion measurements.

Ground-motion displacement is generally obtained either by double time-integration of acceleration data or single time-integration of particle velocity data. This has been the case since no completely satisfactory displacement gauge has been developed. The measurement of particle displacement poses a difficult problem in devising schemes for direct measurement, since a fixed-reference position is required to determine the relative displacement. However, the methods using integration of acceleration and particle velocity data often introduce errors in the calculated displacement. The reference voltage or signal output from a gauge may vary with time (baseline shift) which introduces a cumulative error of unknown value in the displacement obtained by integration. The total signal consists of the real signal superimposed upon the varying reference base. Figure 15 illustrates the different ground displacement results obtained from actual time-integrated accelerometer and particle velocity data. Both gauges were contained in the same instrumentation package. The solid curve shows the results from the acceleration data which have been singly integrated to obtain the particle velocity and doubly integrated to yield displacement. The dotted curve is a record of the velocity measurement and its integration to obtain displacement. These results show the necessity for a method to measure displacement at a point to determine the discrepancy in the measurements.

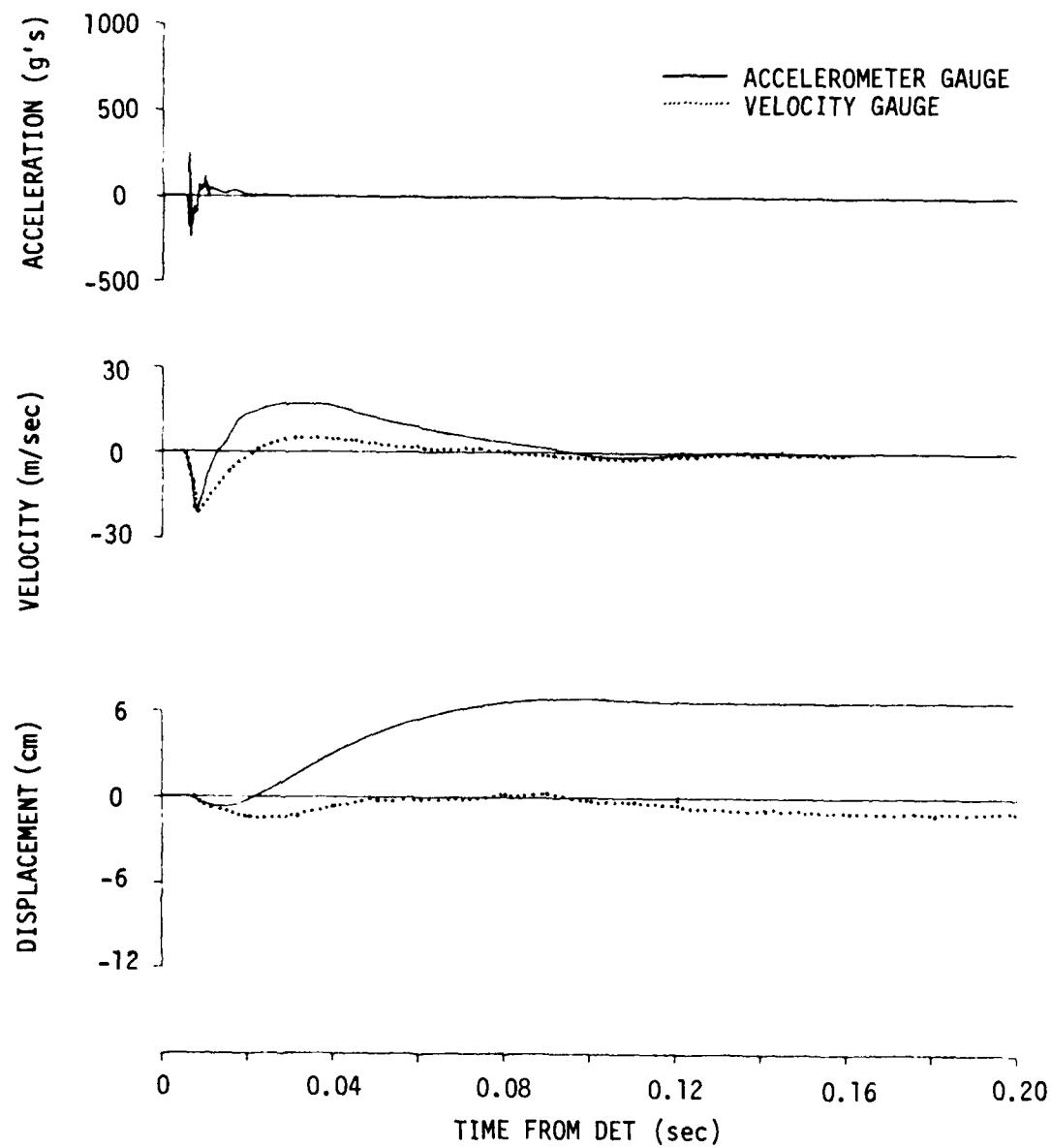


Figure 15. Ground motions vs time, Mineral Rock Gauge 60-18.

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## APPENDIX A. DISPLACEMENT POTENTIAL

From the theory of elasticity, the permanent elastic displacement of the cavity wall is related to the energy released by an explosion given by

$$\frac{4\pi}{3} \frac{Pa^3}{\gamma-1} = \frac{16\pi\mu a^2 d_a}{3(\gamma-1)}$$

where  $p$   $\equiv$  cavity pressure,  
 $\mu$   $\equiv$  shear modulus of medium,  
 $a$   $\equiv$  radius of the cavity,  
 $d_a$   $\equiv$  permanent wall displacement.

With the approximation of incompressible flow,

$$a^2 d_a = r^2 d_r ,$$

the effective energy released in the cavity becomes

$$\frac{4\pi}{3} \frac{Pa^3}{\gamma-1} = \frac{16\pi\mu r^2 d_r}{3(\gamma-1)}$$

The constant  $r^2 d$  is often referred to as the displacement potential. Figure A1 shows the displacement produced by a step-pressure applied to a cavity wall (Ref. 5). The static part of the solution becomes evident at late times. Also shown is the permanent displacement which would be measured by the technique discussed in the report.

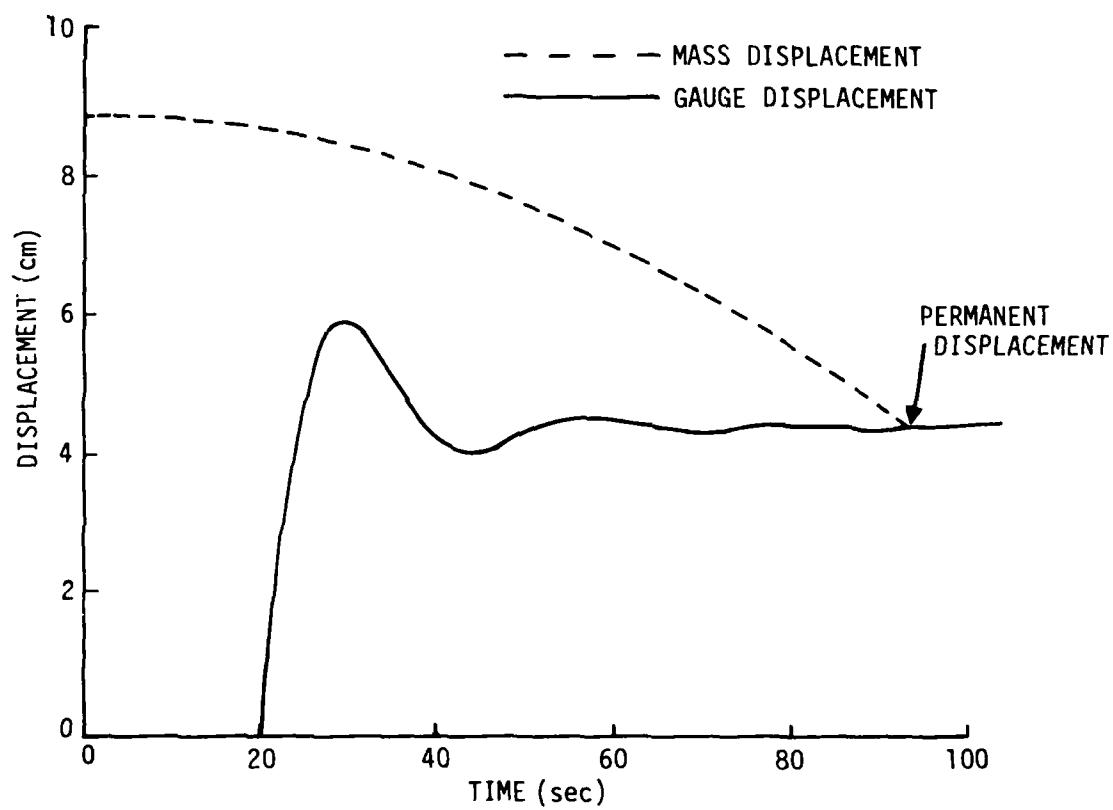


Figure A1. Displacement potential obtained from permanent-displacement measurement.

## APPENDIX B. EFFECTS OF AIR DRAG ON DISPLACEMENT MEASUREMENTS

The relative motion of the drop mass and the air inside the gauge case will apply a drag force opposing the free-fall motion of the weight. In the case of an upward displacement, the flow velocity over the mass will be the ground particle velocity superimposed upon the instantaneous mass velocity. The net force per-unit-area acting on the mass is given by

$$F = 1/2 \rho_a (U_p - U)^2 C_d - mg \quad (B1)$$

where  $m$  = mass per-unit-area of mass,

$C_d$  = drag coefficient,

$U_p$  = gauge case velocity,

$U$  = mass velocity,

$\rho_a$  = air density.

This additional force will affect the time-to-impact and an error in the computed displacement will result. In this Appendix, the error in displacement due to the air drag is estimated.

a. Assumptions--The equations of motion of the mass can be solved using Eq. (B1) when  $U_p$  is known; however, the solutions are quite complicated except for special cases. Several simplifying, conservative assumptions will be made and justified by the results. These assumptions are:

- The particle velocity  $U_p$  is a step function equal to the peak velocity.
- The drag coefficient  $C_d = 1$ .
- The relative velocity  $(U_p + U_o)$  is constant over the motion where  $U_o$  is the maximum free-fall velocity in vacuum,  $U_o = \sqrt{2Sg}$ .
- The mass =  $10 \text{ gm/cm}^2$ .

All of the assumptions are overestimates which will provide an upper bound on the displacement error.

b. Equations of motion--Under these assumptions, the net force  $F$  will be constant over the motion, which will be free-fall with an apparent acceleration  $g'$  obtained from Eq. B1:

$$g' = g \left( 1 - \frac{1/2 \rho_a (U_p - U_o)^2 C_d}{mg} \right) . \quad (B2)$$

The calculated displacement for the measured time-of-fall  $t_m$  then becomes

$$d_c = s - 1/2 g' t_m^2 , \quad (B3)$$

and the displacement in the absence of air drag is given by

$$d = s - 1/2 g t^2 . \quad (B4)$$

Since

$$d_c = U_p t_m \quad (B5)$$

and

$$d = U_p t , \quad (B6)$$

the fractional displacement error can be written

$$\begin{aligned} e &= \frac{d - d_c}{d} \\ &= \frac{t - t_m}{t} \end{aligned} \quad (B7)$$

where  $t$ ,  $t_m$  can be obtained by substituting  $d$ ,  $d_c$  from Eqs. (B5) and (B6) into Eqs. (B3) and (B4) respectively, and solving for  $t$ ,  $t_m$ :

$$t_m = \frac{U_p}{g'} \left( \sqrt{1 + 2Sg' / U_p^2} - 1 \right) \quad (B8)$$

$$t = \frac{U_p}{g} \left( \sqrt{1 + 2Sg / U_p^2} - 1 \right) \quad (B9)$$

where  $U_o + U_p$  in Eq. B2 is given by

$$U_o + U_p = U_p \left( 1 + \sqrt{2gS / U_p^2} \right)$$

to obtain  $g'$ . The parameters in the problem are the gauge size  $S$  and the vertical particle or gauge case velocity  $U_p$ , which will be varied in the displacement error computations.

c. Calculations of displacement error--Table B1 shows values of errors obtained by solving Eqs. (B8) and (B9) using the values of  $S$  and  $U_p$  listed in the table. These values probably cover the range for any practical application of the displacement gauge.

The results show that the drag effects are small for the very conservative approximations used in the computations. There are compensating effects which reduce the error. As the particle velocity increases, the drag also increases; however, the time-to-impact decreases, which tends to reduce the drag impulse on the mass.

TABLE B1. PERCENTAGE ERRORS IN DISPLACEMENT  
DUE TO AIR DRAG

S (cm)	$U_p$ (cm/sec)			
	10	100	1,000	10,000
2	0.01	0.01	0.01	0.01
5	0.03	0.04	0.02	0.02
10	0.07	0.08	0.04	0.03
20	0.14	0.16	0.09	0.07
30	0.20	0.24	0.14	0.10

## APPENDIX C. ERROR DUE TO ROTATIONAL MOTION

Throughout the report in the discussion on the operation of the displacement gauge, it has been assumed that gauge motion was purely translational without rotation. However, a gauge imbedded in a heterogeneous medium or in a region of tectonic stress could undergo a rotation superimposed upon its linear motion.\* In addition, gravity can also introduce rotational motion of the field of flow. In this appendix, the error produced in vertical displacement measurements due to rotational motion is analyzed.

In Figure C1, the initial position of the gauge is shown with the base located along AA. Impact of the mass at P would occur for a vertical displacement  $V$  in the absence of rotation. A rotational motion is assumed for a constant radius of curvature  $r$  through the angle  $\theta$  to the new base position BB. The motion results in a point of impact at  $P'$  with an apparent vertical displacement  $V'$ . Also, the rest point  $P_0$  is displaced horizontally by  $H$ . The error in the vertical displacement with rotation is defined as the fractional difference between the apparent displacement and the true displacement, namely,

$$E = \frac{V' - V}{V}$$

From Figure C1 it can be shown that

$$V' = r \tan \theta$$

$$V = r \sin \theta$$

---

\* Personal communication between the author and G. Rawson (Consultant, RDA).

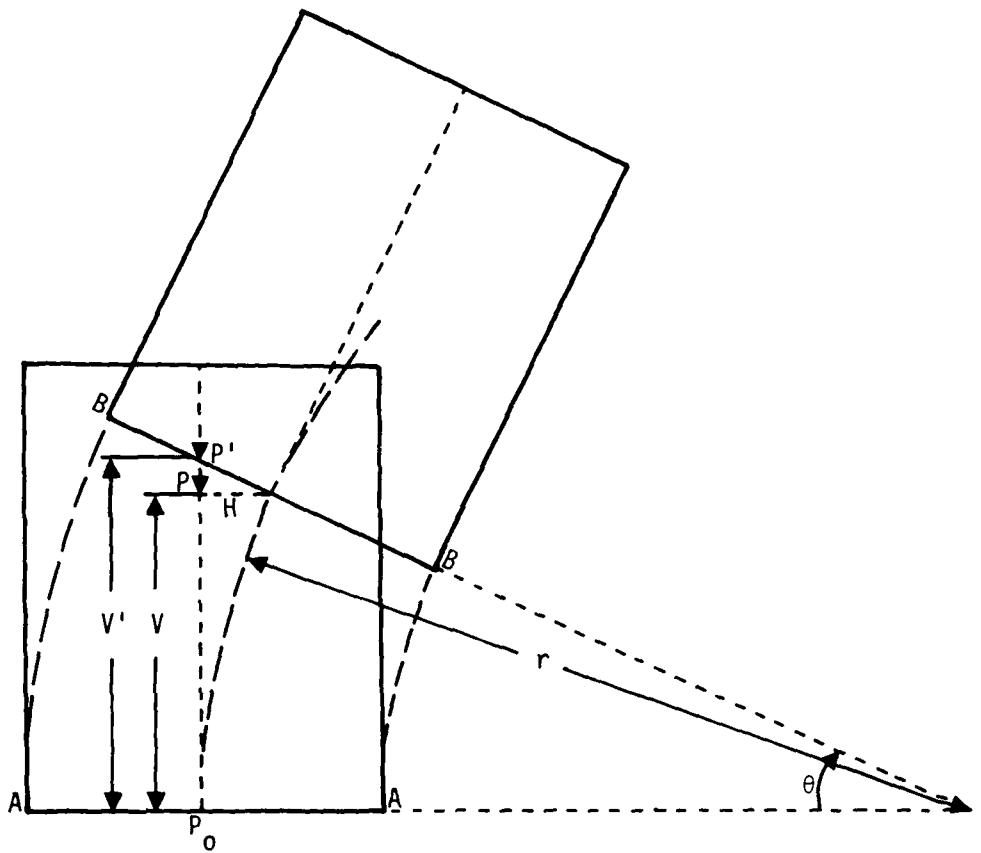


Figure C1. Illustration of effects on ground-motion displacement measurements due to rotational motion.

Hence,

$$\frac{V' - V}{V} = \frac{1 - \cos\theta}{\cos\theta}$$

A plot of the percentage difference between the apparent and true vertical displacements is shown in Figure C2. It is interesting to note that the error is independent of  $r$ , the radius of curvature. Displacements of interest are expected to be small compared to the radius of curvature in rotational motion, hence the angle of rotation would be correspondingly small. Angular motions to 20 deg produce small errors, as noted in Figure C2.

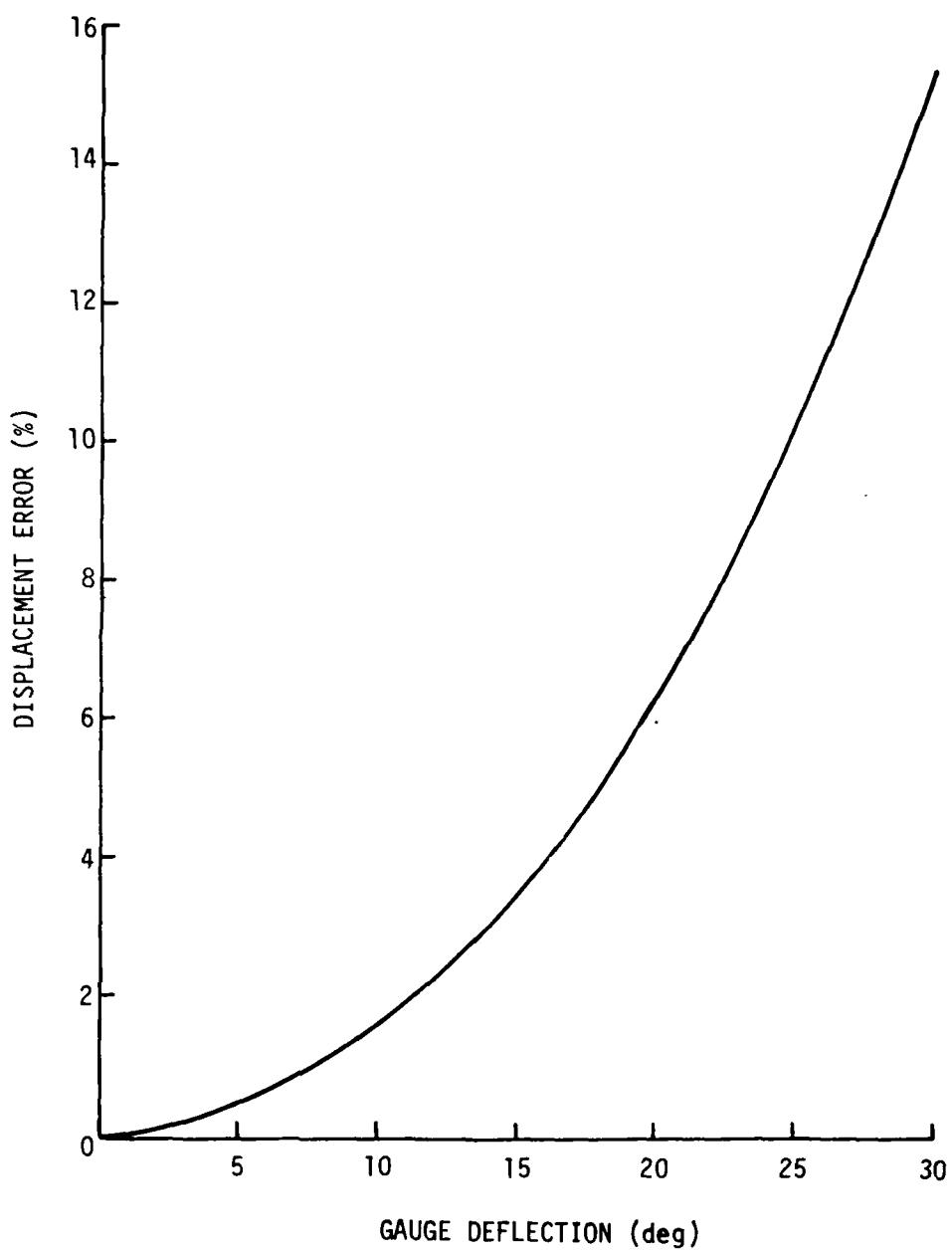


Figure C2. Displacement error due to gauge rotation.

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